

6.S979 Lecture 7

Reminder: Pset 1 due 10/2

Last time:

- Bell-KS theorem
- "contextuality"
- Magic square
→ 2-player game

Today: Generalizations of MS
; a hidden variable theory

1) Linear system game:

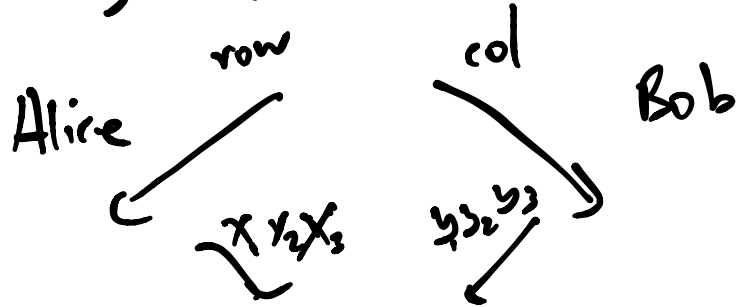
$$\begin{array}{c|c|c} a_1 & a_2 & a_3 = +1 \\ \hline a_4 & a_5 & a_6 \\ \hline a_7 & a_8 & a_9 \\ \hline \vdots & \vdots & \vdots \end{array}$$

Linear equations
over \mathbb{Z}_2

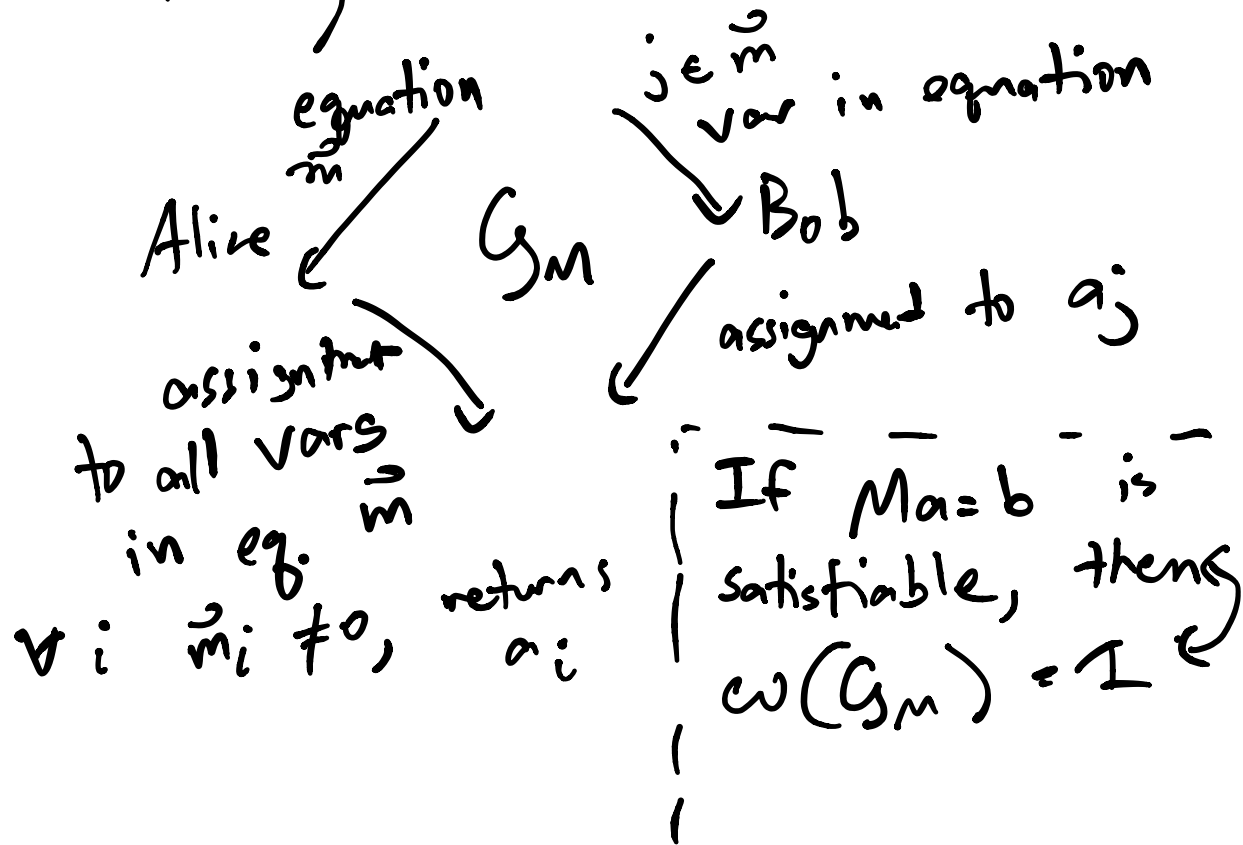
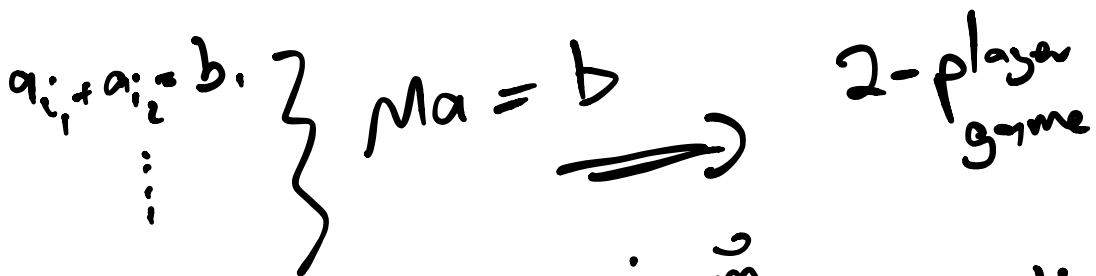
$$a_1 + a_2 + a_3 = 0 \text{ mod } 2$$

$$a_1 + a_4 + a_7 = 0 \pmod 2$$

$$a_3 + a_6 + a_9 = 1 \pmod 2$$



LS game:



If $M_{a=b}$ has a "quantum soln."
then $\omega^*(G_M) = 1$ ↷

Def: A quantum soln. is
a collection of operators
 $\{A_1, \dots, A_n\}$ s.t.

1) $A_i^2 = I \quad \forall i$

2) \forall row $i, \quad \forall j, k$ s.t.
 $M_{ij} \neq 0 \quad \& \quad M_{ik}$ are both nonzero

$$[A_j, A_k] = 0$$

(all vars in an equation should
be simultaneously measurable)

3) \forall row $i, \quad \prod_j A_j^{M_{ij}} = (-1)^{b_i} I$

Can generalize to $\mathbb{Z}_d \pmod{d}$

1) $A_i^d = I$, 2) $\prod_j A_j^{m_{ij}} = \omega^b I$

eigenvalues are d -th roots of unity
 "generalized observables"

$\omega = e^{2\pi i/d}$

Q. sol \Rightarrow val. \perp Q. strat

$$\frac{1}{\sqrt{d}} \sum_i |i\rangle |i\rangle$$

$$A_i \quad A_i^*$$

consistent w/ prob I

$\{A_i\}$ q. sol. \Rightarrow satisfy eq. w/ prob I

val. \perp Q. strat \Rightarrow Alice's observables are a q. sol

Def. Solution group

Generators: $\{a_1, \dots, a_n, J\}$

1. $a_i^d = J^d = e \leftarrow$ identity

2. $a_i J = J a_i \quad \forall i$

3. $\forall i, j$ in same equation,
 $a_i a_j = a_j a_i$

4. $J^{-b_i} \prod_j a_j^{M_{ij}} = e \quad \forall i$

A quantum sol. is a unitary rep
of sol. group where

$$J \mapsto \text{co. I}$$

Consequence:

If from sol. group relations,
can prove that $J = e$, then
 \nexists quantum sol.

[Cleve, Liu, Slofstra]
If d prime and $J \neq e$
then \exists q. sol.

MS mod d (d odd prime)

$$a_1 \ a_2 \ a_3 = 0 \text{ mod } d$$

$$a_4 \ a_5 \ a_6 = 0 \text{ mod } d$$

$$a_7 \ a_8 \ a_9 = 0 \text{ mod } d$$

$$0 \quad 0 \quad 1 \text{ mod } d$$

No classical solution

Sol. group relations
 $\Rightarrow J^2 = e = J^d$

$$\Rightarrow J = e$$

\Rightarrow No q. sol

Generalized Paulis

$$X = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \quad Z = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$$
$$XZ = -ZX \quad X^2 = Z^2 = I$$

$$X = \sum_i |i+1\rangle \langle i|$$

$$Z = \sum_i \omega^i |i\rangle \langle i|$$

$$X^d = Z^d = I$$

$$ZX = \sum_i \omega^{i+1} |i+1\rangle \langle i|$$

$$= \omega \left(\sum_i \omega^i |i+1\rangle \langle i| \right)$$

$$= \omega \cdot XZ$$

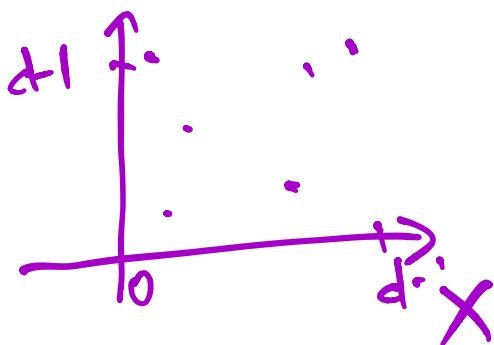
$MS \rightarrow$ "tests anticommutation"
 $A_1 A_5 = -A_5 A_1$

$MS_d \rightarrow$ tests phase commutation
 $a_1 a_5 = \zeta a_5 a_1$

$$\Rightarrow \zeta^2 = 1$$

There is a hidden variable
 model for mod d $X \hat{=} Z$
 operators acting on a big class
 of states

"Quasiprobability distribution
 over phase space"
 $x \hat{=} z$ values



$$W(x, z) \neq 0$$

$$\sum_z W(x, z) = \text{Pr}_\psi[x]$$

$$\sum_x W_\psi(x, z) = \text{Pr}_\psi[z]$$

Wigner function W
for states & measurements

$$W_\psi(x, z) = \frac{1}{d} \langle \psi | A(x, z) | \psi \rangle$$

$$W_{M_k}(x, z) = \text{tr}(A(x, z) M_k)$$

$$\begin{aligned} \text{Pr}[k] &= \langle \psi | M_k | \psi \rangle \\ &= \sum_{x, z} W_\psi(x, z) \cdot W_{M_k}(x, z) \end{aligned}$$

$$= \mathbb{E}_{(x, z) \sim W_\psi} W_{M_k}(x, z)$$

hidden state \nearrow (x, z) sampled
from W_ψ

Fact: 1) For any mod d measurement
 X or Z

$$W_{M_R} \geq 0$$

2) \exists states ψ for which
 $W_{\psi} \geq 0$

Thm: [Qassim & Wallman]

Can map any q. sol. built
out of mod d X, Z , & products
to a classical sol.

Pf: Fix the hidden state to
 $(x, z) = (0, 0)$

$$W_{\psi}(x, z) = \delta_{x=0, z=0}$$

$$W_{M_k}(x, z) = \text{tr}(A(x, z) M_k)$$

$$A(0, 0) = \sum_i | -i \rangle \langle i |$$

$$= \sum_i | i + 2^{-1}d \rangle \langle i |$$

for any observable $\omega^k Z^p X^q$

Wigner function gives

det. outcome $v(\omega^k Z^p X^q)$

$$= k + 2^{-1} p q$$

↑
multiplicative inverse mod d

$\{A_1, \dots, A_n\}$ q. sol.

$\Rightarrow \{v(A_1), \dots, v(A_n)\}$ d. sol.

$$v(AB) = v(A)v(B)$$

if A & B commute

$$A = Z^p X^q, \quad B = Z^r X^s$$

$$ZX = \omega XZ \quad [A, B] = 0 \quad \text{if}$$

$$ps - qr = 0$$

$$r(AB) = r(Z^p X^q Z^r X^s)$$

$$= r(\omega^{-rq} Z^{p+r} X^{q+s})$$

$$= -rq + 2 \cdot (p+r)(q+s)$$

$$= -1 \cdot r \cdot q + 2(p \cdot q + r \cdot s)$$

$$+ 2(p \cdot s + r \cdot q)$$

$$= 2(p \cdot q) + 2(r \cdot s) + 2(p \cdot s - r \cdot q) = 0$$

$$= r(A) + r(B) + 0$$

Takeaway:

A HVN can describe a
subtheory of QM with

— Non-compatible measurements

— Entanglement

" Stabilizer states & Clifford " circuits "

Caveat:

— mod 2 vs mod d

— stabilizers & cliffords are computationally
simulable mod 2, but no HVN

Linear systems games and self-testing

Thm [Coladangelo Stark]:

Suppose soln. group has a unique irrep σ where

$$J \xrightarrow{\sigma} \text{co } I$$

Then you have a robust self-test for σ

$$(V_A \otimes V_B)(A_i \otimes I |\psi\rangle)$$

$$\approx (\sigma(a_i) \otimes I |\hat{\psi}\rangle | \text{junk} \rangle$$

"Self-test for measurements"